

# Sintering behavior and microwave dielectric properties of CaTi<sub>1-x</sub>(Nb<sub>1/2</sub>Al<sub>1/2</sub>)<sub>x</sub>O<sub>3</sub>



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## ABSTRACT

CaTi<sub>1-x</sub>(Nb<sub>1/2</sub>Al<sub>1/2</sub>)<sub>x</sub>O<sub>3</sub> with x=0.1–0.5 ceramics were processed through solid state sintering. X-rays diffraction (XRD) patterns of the compositions showed that the samples have orthorhombic crystal structure with symmetry (Pbnm). The symmetry was further confirmed using Raman spectroscopy. A total of 13 Raman modes were detected, which were in agreement with the XRD results. Microstructure analysis of the samples showed porosity in the samples, presumably due to the substitution of Al, having high melting point. As the concentration of Al and Nb increased, relative permittivity ( $\epsilon_r$ ), quality factor (Q×f<sub>0</sub>) and temperature coefficient of resonance frequency decreased. Optimum microwave dielectric properties were achieved for the composition x=0.5 sintered at 1650 °C for 8 h i.e.,  $\epsilon_r \sim 27.09$ , Q×f<sub>0</sub> ~17378 GHz, and  $\tau_r \sim -2.5$  ppm/°C.

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# 1. Introduction

Dielectric ceramics are widely used in different communication systems such as cellular phones, satellites, base stations, military radar, navigation system, ITS, automation, filters, oscillators etc. For microwave (MW) applications, appropriate relative permittivity  $(\varepsilon_r)$ , low tan  $\delta$  or high Q×f<sub>0</sub> and near zero temperature coefficient of resonance frequency  $(\tau_f)$  are preferred [1–3]. For MW applications, perovskites are of special interest due to their interesting physical and chemical properties. Perovskite structure CaTiO3 was discovered by German mineralogist, Gustav Rose (in 1839) who he named as perovskite in the recognition of a Russian mineralogist L.A. Perovski. CaTiO<sub>3</sub> has orthorhombic structure and exhibit high  $\varepsilon_r = 162$ , reasonable Q×f<sub>o</sub> (~8354 GHz) but its high  $\tau_f$  (+803.7 ppm/°C) renders its use in practical applications [4]. To tune  $\tau_f$  of CaTiO<sub>3</sub>, A- and/or Bsite or both sites have been modified. Tao et al. [5] investigated  $[Ca_{1-x}(Li_{1/2}Nd_{1/2})]TiO_3$ , while substituting  $Ca^{2+}$  by  $(Li_{1/2}Nd_{1/2})^{2+}$ at the B-site and obtained optimal MW dielectric properties;  $\epsilon_r = 112.6$ ,  $Q \times f_o = 4480$  GHz and  $\tau_f = +8.2$  ppm/°C at x = 0.

KEYWORDS

Perovskite CaTiO<sub>3</sub> Microwave dielectric properties

(1-x)CaTiO<sub>3</sub>-xCa(Ta<sub>1/2</sub>Ga<sub>1/2</sub>)O<sub>3</sub> possess  $\varepsilon_r = 47$ ,  $Q \times f_o = 26,630$  GHz and  $\tau_f \sim 2.64$  ppm/°C for the x = 0.5 [6]. The effect of La substitution for Ca in CaTiO<sub>3</sub> with general formula Ca<sub>1-x</sub>La<sub>2x/3</sub>O<sub>3</sub> was exhibited  $Q \times f_o = 17600$  GHz,  $\varepsilon_r = 109$ ,  $\tau_f = +213$  ppm/°C at x = 0.4 [7]. Ca<sub>0.66</sub>La<sub>0.387</sub>Ti<sub>0.88</sub>O<sub>3</sub> possess high  $Q \times f_o = 13407$  GHz and  $\varepsilon_r = 71.5$  but a poor  $\tau_f = +136.4$  ppm/°C [8]. Feteira et al. [9] observed a decrease in  $\varepsilon_r$  from 170 to 12, while an increase in  $Q \times f_o$  from 10000 to 12000 GHz in compositions Ca<sub>1-x</sub>Y<sub>x</sub>Ti<sub>1-x</sub>Al<sub>x</sub>O<sub>3</sub> (x = 0.1). 1-x(CaTiO<sub>3</sub>)-x(Li<sub>0.5</sub>La<sub>0.5</sub>)TiO<sub>3</sub> compositions (x = 0.2–0.8) exhibit  $\varepsilon_r = 245$ ,  $Q \times f_o = 2750$  GHz, and  $\tau_f = 0.75$  ppm/°C at x = 0.6 [10]. A niobate-based system, Ca(Fe<sub>1/2</sub>Nb<sub>1/2</sub>)O<sub>3</sub> has  $\varepsilon_r = 40$ , the  $Q \times f_o = 20000$  GHz, and  $\tau_f = -76$  ppm/°C [11]. Similarly, Nb- and Gamodified CaTiO<sub>3</sub> sintered temperature 1350–1475 °C for 4 h possess a good set of MW properties i.e.,  $\varepsilon_r = 52$ ,  $Q \times f_o = 23,595$  GHz and  $\tau_f = -4.9$  ppm/°C [12].

Another system, CaTiO<sub>3</sub>-Ca(Zn<sub>1/3</sub>Nb<sub>2/3</sub>)O<sub>3</sub> was investigated for MW dielectric properties which revealed  $\epsilon_r = 51$ , Q×f<sub>o</sub> = 10860 GHz and  $\tau_f \sim$ -6 ppm/°C [13]. The effect of niobium doping in CaTiO<sub>3</sub> is very prominent at the B-site. Therefore, we have doped Ti<sup>+4</sup> by (Al<sub>1/2</sub>Nb<sub>1/2</sub>)<sup>4+</sup>

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to maintain charge balance, and its effect on the microwave dielectric properties was investigated.

#### 2. Experimental techniques

 $CaTi_{1-x}(Nb_{1/2}Al_{1/2})_{x}O_{3}$  (x = 0.1, 0.2, 0.3, 0.4, and 0.5) ceramic solid solution was prepared via solid state sintering route. For accurate weighing, the raw materials CaCO<sub>3</sub>, TiO<sub>2</sub>, Nb<sub>2</sub>O<sub>5</sub> and Al<sub>2</sub>O<sub>3</sub> were dried to remove the moistures and hydroxide (if any). The dried powders were weighted and then ball-milled in isopropanol for 24 h using Ytoughened zirconia balls as a grinding media. After mixing, the powders were calcined at 950 °C for 6 h in alumina crucible at a heating/cooling rate of 10 °C/min. The calcined powders were remilled and then pressed into pellets at a pressure of 100 MPa in cylindrical stainless-steel die with a 10 mm in diameter using a uniaxial pellet press. Green pellets were sintered in the temperature range of 1450-1650 °C for 8 h in a muffle furnace to achieve maximum density. The density of the sintered discs was measured using a high precession densitometer (DM 500) using the Archimedes method. For phase identification of the samples, x-ray diffractometer (PANalytical X'pert Pro) was used. The diffraction patterns were recorded at room temperature (RT) using Cu-K $\alpha$  radiations ( $\lambda = 1.5405$  Å) in 2 $\theta$  range from 10-70°, at a small step size 0.05° and scanning time of 3 sec/step. Raman spectra were recorded at RT using a Renishaw Raman microscope with a diode 514 nm excitation Ar laser. The exposure time was 10 s for each sample. Surface morphology was studied using scanning electron microscope JEOL (JSM-6460LV). Before SEM analysis, pellets were polished, thermally etched and carbon coated. MW dielectric properties ( $\epsilon_r$ , Q×f<sub>o</sub>, and  $\tau_f$ ) of the sintered sample were measured using an Advantest-R3767CH vector network analyzer.  $\tau_f$  was measured in temperature range from RT to 80 °C.

## 3. Results and discussion

X-Ray Diffraction (XRD) patterns of the sintered  $CaTi_{1-x}(Nb_{1/2}Al_{1/2})_xO_3$  (CTNA) for x = 0.1–0.5 ceramics are shown in Fig. 1. The Braggs reflections matched PDF # 86-1393 for all samples which confirm the

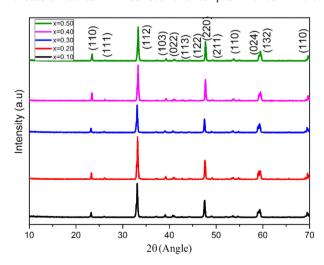


Fig. 1. XRD patterns of CTNA samples, sintered at their optimum sintering temperature.

orthorhombic structure having space group Pbnm (No. 62). XRD peaks slightly shifted to higher 2 $\theta$  angles with increase in x, which may be attributed to smaller ionic radius of  $[Al_{1/2}Nb_{1/2}]^{4+}$  (r = 0.585 Å) in comparison to Ti<sup>4+</sup> (r = 0.605 Å) [14]. Thus, the unit cell volume decreased with increase in x.

Raman spectroscopy analysis was also carried out to further confirm the symmetry of the samples, Fig. 2. A total of 13 Raman modes were detected (Table 1) which are in agreement with previous studies [12, 15. 16]. The band near 637 cm<sup>-1</sup> for the sample with x = 0.1 can be allocated to the Ti-O symmetric stretching mode [17], which disappeared as the Ti4<sup>+</sup> is replaced by the dopant Al<sup>3+</sup> and Nb<sup>5+</sup>. The intensity of the band near ~804 cm<sup>-1</sup> increased with increasing concentration which is assigned to the rotation of BO<sub>6</sub> octahedra [12]. As the ionic radius of Al<sup>3+</sup>, Nb<sup>5+</sup> and Ti<sup>+4</sup> is different; therefore, their bands will appear at slightly different frequencies. Due to nonequivalence their vibration may become Raman active modes as stated by Zheng et al. [16]. The band near 838 cm<sup>-1</sup> starts appearing as x is increased which is due to multiple cations at the B-site, as reported for CaTiO<sub>3</sub> based compounds having more than one cations at the Ti-site, known as A1g mode [18]. Raman bands at ~244, 471, and 804 cm<sup>-1</sup> may be associated with the asymmetric and symmetric breathing of octahedra [19]. The peaks at 499 cm<sup>-1</sup> may be attributed to Ti-O<sub>3</sub> torsional mode which shifted to 539 cm<sup>-1</sup> which is due to the fact that Ti-O-Ti bond length is larger than Al-O-Nb [20]. Cockayne et al. [21] suggested that it may also be due to cation ordering of Al-O-Nb bond. The bands at ~121, 149, 178, 224, 277, and 337 cm<sup>-1</sup> are assigned to O-Ti-O bending and 244 cm<sup>-1</sup> may be attributed the Ca-TiO<sub>3</sub> lattice mode [12, 16].

SEM micrographs of the thermally etched and polished CTNA samples sintered at 1650 °C for 8 h are shown in Fig. 3. Some voids were observed in the SEM images, thus leading to a low density, consistent with the apparent density of the sample. A granular like connected grains can be seen in the microstructure of sample x = 0.1. For sample x = 0.2 and 0.3, hexagon like grains were formed and the grain size increased. In comparison to sample x = 0.1, the x = 0.2 sample has less voids and hence greater density. With further increase in x, a clear increase in the grain size as well as a slight decrease in density was observed. For sample x = 0.5, well-connected grains with some voids were observed; however, the apparent density was still lower due to these larger voids.

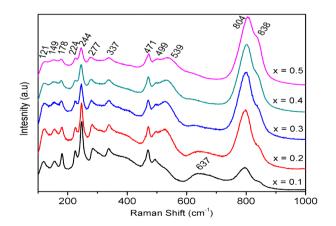


Fig. 2. Raman spectra of CTNA samples.

Band	1	2	3	4	5	6	7	8	9	10	11	12	13	Ref.
CTNA	121	149	178	224	244	277	337	471	499	539	637	804	838	Present work
CTNG	123	153	180	225	249	279	345	461	485	543	641	792	817	[12]
YCT-NA	-	163	183	227	247	288	339	470	494	509	641	-	825	[16]

Table 1. Comparison of the Raman bands of CTNA with literature.

MW dielectric properties of CTNA samples are given in Table. 2. At x = 0.1, the sample possess  $\varepsilon_r = 55$  which increased with increase in x from 0.1 to 0.2. The only possible reason for the increase in relative permittivity may be the lower density of x = 0.1 sample in comparison to x = 0.2 which is also evident from the micrograph and observed

density of these samples. With further increase in x from 0.2 to 0.5,  $\varepsilon_r$  gradually decreased. Also,  $\varepsilon_r$  strongly depends on ionic polarizability and molar cell volume. In the present case, both ionic polarizabilities per unit volume decreased due to the substitution of  $(Al_{1/2}Nb_{1/2})^{4+}$  for Ti<sup>4+</sup> [14]. On the other hand, the Q×f<sub>o</sub> showed similar trend like  $\varepsilon_r$  for

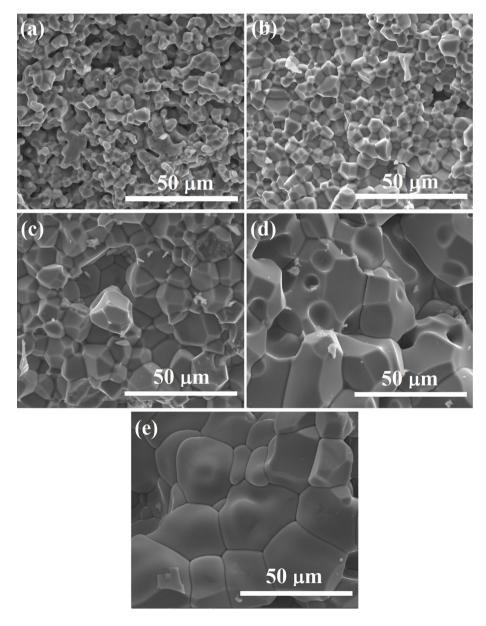


Fig. 3. SEM images of CTNA samples, sintered at 1650 °C for 8 h; x = a) 0.1, b) 0.2, c) 0.3, d) 0.4, and e) 0.5.

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Sample (x)	Sintering temperature (°C)	Sintering time (h)	Density (g/cm <sup>3</sup> )	٤r	Q×f <sub>0</sub> (GHz)	τ <sub>f</sub> (ppm/°C)
0.1	1650	8	2.88	55.11	7754	+456.6
0.2	1650	8	3.548	75.08	11487	+274.3
0.3	1650	8	3.498	57.80	11612	+137.13
0.4	1650	8	3.35	45.39	15390	+35.92
0.5	1650	8	2.903	27.09	17378	-2.5

Table 2. MW dielectric properties of CTNA.

the samples x = 0.1 and 0.2 because  $Q \times f_o$  is also strongly dependent on density [22]. As the porosity decreased, that quality factor increased.

With an increase in x, the quality factor significantly enhanced which may be due to a decrease in lattice harmonic vibrations due to substitution of lower ionic polarizability and smaller ionic radii of  $Al^{3+}$  in comparison to  $Ti^{4+}$  [23]. Similarly, grain boundaries per unit volume decrease, which limit the conduction and hence improves the quality factor [6]. Generally,  $\epsilon_r$  and  $\tau_f$  have a linear relationship because  $\tau_f$  is also influenced by the temperature coefficient of relative permittivity [24, 25].  $\tau_f$  value decreased from +456.6 to -2.5 ppm/°C which may be attributed to a decrease in tilting angle of BO<sub>6</sub> octahedra [26]. The tolerance factor of CTNA is 0.972 which is higher than CT (0.965); therefore, an anti-phase tilting scheme is expected [27]. Mostly, in this tiling scheme, perovskites exhibit a near zero temperature coefficient of dielectric constant and hence temperature coefficient of the resonance frequency.

# 4. Conclusions

In the present work, CaTi<sub>1-x</sub>(Nb<sub>1/2</sub>Al<sub>1/2</sub>)xO<sub>3</sub> (x = 0.1 to 0.5) were prepared via solid state sintering route. XRD results and Raman results confirmed the formation of single-phase orthorhombic crystal structures (Pbnm). The effect of sintering temperature on the microwave dielectric properties was also investigated. The samples were found to become harder in density; therefore, sintering time was increased. SEM results showed some voids in the microstructures which were associated with the lower density of the samples.  $\varepsilon_r$ ,  $\tau_f$  decreased and Q×f<sub>0</sub> was observed to increase as x from x = 0.2 to x = 0.5. The optimum set of MW dielectric properties was observed for sample x = 0.5 i.e.,  $\varepsilon_r \sim 27$ , Q×f<sub>0</sub>  $\sim 17378$  GHz and high  $\tau_f \sim 2.5$  ppm/°C when sintered temperature 1650 °C for 8 h which may be suitable for use in practical applications.

# **CRediT** authorship contribution statement

Akbar Khan: Methodology, Writing – original draft. Asif Ali: Supervision, Writing – review & editing. Izaz Khan: Investigation, Writing – review & editing.

#### Data availability

The data underlying this article will be shared on reasonable request to the corresponding author.

#### **Declaration of competing interest**

The authors declare no competing interests.

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